

Optimization-Based Mesh Correction

M. D'Elia, D. Ridzal, <u>K. Peterson</u>, P. Bochev Sandia National Laboratories <u>M. Shashkov</u> Los Alamos National Laboratory

> MultiMat September 9, 2015

SAND 2015-7334C



Sandia National Laboratories is a multi-program laboratory managed and operated by Sandia Corporation, a wholly owned subsidiary of Lockheed Martin Corporation, for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-AC04-94AL85000.



Divergence Free Lagrangian Motion

• Given cell κ at initial time t^0





Divergence Free Lagrangian Motion

- Given cell κ at initial time t^0
- Compute nodal displacement from velocity field u





Divergence Free Lagrangian Motion

- Given cell κ at initial time t^0
- Compute nodal displacement from velocity field u
- Updated cell k̂(t¹) has both temporal and spatial errors

Violation of volume preservation $\frac{d}{dt}\int_{\kappa(t)}dV\neq\mathbf{0}$





Divergence Free Lagrangian Motion

- Given cell κ at initial time t^0
- Compute nodal displacement from velocity field u
- Updated cell k̂(t¹) has both temporal and spatial errors

Violation of volume preservation $\frac{d}{dt}\int_{\kappa(t)}dV\neq\mathbf{0}$

Consider $\rho = const$

Let cell mass
$$M_{\kappa}(t) = \int_{\kappa(t)} \rho dV$$
 and cell density $\rho_{\kappa} = \frac{M_{\kappa}(t)}{|\kappa(t)|}$,
where $|\kappa(t)| = \int_{\kappa(t)} dV$

$$\rho_{\kappa}(t^{1}) = \frac{M_{\kappa}(t^{1})}{|\hat{\kappa}(t^{1})|} \neq \frac{M_{\kappa}(t^{0})}{|\kappa(t^{0})|} = \rho_{\kappa}(t^{0})$$

Cannot maintain a constant density!







Geometric Conservation Law (GCL)

$$\frac{d}{dt}\int_{\kappa_i(t)} dV = \int_{\partial \kappa_i(t)} \mathbf{u} \cdot \mathbf{n} \, ds$$

Some recent work:

Use more Lagrangian points

Enforces GCL approximately

Lauritzen, Nair, Ullrich (2010), A conservative semi-Lagrangian multi-tracer transport scheme on the cubed-sphere grid, *JCP*.

Heuristic mesh adjustment procedure

No theoretical assurance of completion

Arbogast and Huang (2006), A fully mass and volume conserving implementation of a characteristic method for transport problems, *SISC*.

Monge-Ampére trajectory correction

Accuracy of the MAE scheme determines accuracy of GCL approximation

Cossette, Smolarkiewicz, Charbonneau (2014), The Monge-Ampere trajectory correction for semi-Lagrangian schemes, *JCP*.





- Adjusted point to remain fixed at this stage.
- Points adjusted simultaneously in the direction of the characteristic.
- × Points adjusted "sideways" to the flow.

$$egin{aligned} & ilde{\mathbf{p}}_{ij}^{corr} = ilde{\mathbf{p}}_{ij} + (t - t_n)
abla \phi; \ & ext{det} \ rac{\partial \mathbf{p}_{ij}}{\partial x} = 1 \end{aligned}$$



Optimization-Based Solution

Given a *source mesh* $\widetilde{K}(\Omega)$, and *desired cell volumes* $c_0 \in \mathbb{R}^m$ such that

$$\sum_{i=1}^m c_{0,i} = |\Omega|$$
 and $c_{o,i} > 0 orall i = 1,...,m$

Find a *volume compliant mesh* $K(\Omega)$ such that

- (1) $K(\Omega)$ has the same connectivity as the source mesh
- 2 The volumes of its cells match the volumes prescribed in c_0
- **3** Every cell $\kappa_i \in K(\Omega)$ is valid or convex
- **Output** Boundary points in $K(\Omega)$ correspond to boundary points in $\widetilde{K}(\Omega)$



 $p_{i,3} = (x_{i,3}, y_{i,3})$

Requirements for Quadrilateral Cells

Oriented volume of quad cell:

$$c_i(K(\Omega)) = \frac{1}{2} ((x_{i,1} - x_{i,3})(y_{i,2} - y_{i,4}) + (x_{i,2} - x_{i,4})(y_{i,3} - y_{i,1}))$$

Partitioning of quad into triangles:

$$(a_r, b_r, c_r) = \begin{cases} (1, 2, 4) & r = 1\\ (2, 3, 4) & r = 2\\ (1, 3, 4) & r = 3\\ (1, 2, 3) & r = 4. \end{cases}$$

Oriented volume of triangle cell:



 $p_{i,4} = (x_{i,4}, y_{i,4})_{\bullet}$

$$t_i^r(K(\Omega)) = \frac{1}{2}(x_{i,a_r}(y_{i,c_r} - y_{i,b_r}) - x_{i,b_r}(y_{i,a_r} - y_{i,c_r}) - x_{i,c_r}(y_{i,b_r} - y_{i,a_r})).$$

Convexity indicator for a quad cell:

$$\mathcal{I}_i(K(\Omega)) = \left\{egin{array}{cc} 1 & ext{if } orall t_i^r \in \kappa_i, |t_i^r| > 0 \ 0 & ext{otherwise} \end{array}
ight.$$



Optimization Problem

Objective:

Mesh distance
$$J_0(\boldsymbol{p}) = \frac{1}{2}d(K(\Omega), \tilde{K}(\Omega))^2 = |\boldsymbol{p} - \tilde{\boldsymbol{p}}|_{l^2}^2$$

Constraints:

- (1) Volume equality
- (2) Cell convexity
- $\forall \kappa_i \in K(\Omega), |\kappa_i| = c_{0,i}$
- $\forall \kappa_i \in K(\Omega), \forall t_i^r \in \kappa_i, |t_i^r| > 0$
- (3) Boundary compliance $\forall p_i \in \partial \Omega, \gamma(p_i) = 0$

Nonlinear programming problem (NLP)

 $p* = \arg \min \{J_0(p) \text{ subject to } (1), (2), \text{ and } (3)\}$



Simplified Formulation

For polygonal domains

- boundary compliance can be subsumed in the volume constraint
- convexity can be enforced weakly by logarithmic barrier functions

Objective:

Mesh distance - log barrier
$$J(p) = J_0(p) - \beta \sum_{i=1}^m \sum_{r=1}^4 \log t_i^r(p)$$

 $J_0(p) = |p - \tilde{p}|_{l^2}^2$

Constraints:

(1) Volume equality $\forall \kappa_i \in K(\Omega), |\kappa_i| = c_{0,i}$

Simplified NLP

$$oldsymbol{p}^* = rg \min ig \{ J(oldsymbol{p}) ext{ subject to (1)} ig \}$$



Scalable Optimization Algorithm

Based on the **inexact trust region** sequential programming (SQP) method with key properties:

- Fast local convergence
- Use of very coarse iterative solvers
- Handles rank-deficient constraints

Linear systems for an optimization iterate p^k are of the form

$$\left(egin{array}{ccc} I &
abla C(oldsymbol{p}^k)^T \
abla C(oldsymbol{p}^k) & 0 \end{array}
ight) \left(egin{array}{ccc} oldsymbol{v}^1 \ oldsymbol{v}^2 \end{array}
ight) = \left(egin{array}{ccc} oldsymbol{b}^1 \ oldsymbol{b}^2 \end{array}
ight),$$

where $C(p^k)$ is a nonlinear vector function of coordinates.

Preconditioner

$$\pi^k = \left(egin{array}{ccc} I & 0 \ 0 & \left(
abla oldsymbol{C}(oldsymbol{p}^k)
abla oldsymbol{C}(oldsymbol{p}^k)^T + arepsilon I
ight)^{-1} \end{array}
ight)$$

- $\epsilon > 0$ small parameter $10^{-8}h$
- $\nabla C(p^k) \nabla C(p^k)^T$ formed explicitly
- Smoothed aggregation AMG used for inverse

Heinkenschloss, Ridzal (2014), A matrix-free trust-region SQP method for equality constrained optimization SIOPT.



Scalablility Test

To challenge the algorithm we test performance as follows:

- Start with uniform $n \times n$ mesh and advance to final time using velocity field
- Apply algorithm to the deformed mesh at the final time using initial mesh volumes

n	SQP	CG	GMRES tot.	GMRES av.	CPU	% ML time
64	5	15	101	2.8	2.475	66
128	4	9	106	4.1	8.799	78
256	5	5	130	5.0	45.733	83
512	6	1	100	3.8	184.446	83
Gauss-Newton Hessian						
n	SQP	CG	GMRES tot.	GMRES av.	CPU	% ML time
64	5	5	64	2.5	1.666	63
128	4	4	79	3.8	6.466	82
256	5	5	126	4.8	43.241	86
512	6	6	100	3.8	183.697	86

Analytic Hessian

Almost constant GMRES iterations

• CPU per SQP iteration scales linearly with problem

Computational time dominated by AMG (ML) preconditioner

• The algorithm inherits its scalability from the AMG solver



Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$





Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$













Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$













Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$

8x8 mesh, T = 8, CFL = 2, $\beta = 0$, forward Euler for trajectories.



Uncorrected









Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$

8x8 mesh, T = 8, CFL = 2, $\beta = 0$, forward Euler for trajectories.



Uncorrected



Optimized





Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$

8x8 mesh, T = 8, CFL = 2, $\beta = 0$, forward Euler for trajectories.



Uncorrected



Optimized





Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$





Uncorrected









Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$

8x8 mesh, T = 8, CFL = 2, $\beta = 0$, forward Euler for trajectories.





Uncorrected









Swirling velocity field:
$$\mathbf{u}(\mathbf{x}, t) = \begin{pmatrix} \cos\left(\frac{t\pi}{T}\right)\sin(\pi x)^2\sin(2\pi y) \\ -\cos\left(\frac{t\pi}{T}\right)\sin(\pi y)^2\sin(2\pi x) \end{pmatrix}$$





Uncorrected









Improvements in Mesh Geometry

We observe significant improvements in the geometry of the corrected mesh:

• The shapes of the corrected cells are closer to the exact Lagrangian shapes

Exact



Uncorrected









Improvements in Mesh Geometry

We observe significant improvements in the geometry of the corrected mesh:

- The shapes of the corrected cells are closer to the exact Lagrangian shapes
- The barycenters of the corrected cells are closer to the exact barycenters











Improvements in Mesh Geometry

We observe significant improvements in the geometry of the corrected mesh:

- The shapes of the corrected cells are closer to the exact Lagrangian shapes
- The barycenters of the corrected cells are closer to the exact barycenters
- The trajectories of the corrected cells are closer to the exact Lagrangian trajectories





 $\beta = 0$





 $\beta = 0$





 $\beta = 2.0 imes 10^{-6}$





 $eta=2.0 imes10^{-5}$





Want to solve:
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = \mathbf{0}$$

Given cell volume $c_i = \int_{\kappa_i} dV$, cell mass $m_i = \int_{\kappa_i} \rho(\mathbf{x}, t) dV$, and cell average density $\rho_i = m_i/c_i$ at time t



- $lacksymbol{1}$ Define Lagrangian departure cells: $c_i
 ightarrow ilde c_i$
- 2) Remap from fixed grid to departure grid: $ho_i o ilde
 ho_i, \, ilde m_i = ilde
 ho_i ilde c_i$
- 3 Lagrangian update: $m_i(t + \Delta t) = \tilde{m}_i$, $\rho_i(t + \Delta t) = m_i(t + \Delta t)/c_i$

Sandia

Application to Semi-Lagrangian Transport

Want to solve:
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = \mathbf{0}$$

Given cell volume $c_i = \int_{\kappa_i} dV$, cell mass $m_i = \int_{\kappa_i} \rho(\mathbf{x}, t) dV$, and cell average density $\rho_i = m_i/c_i$ at time t



-) Define Lagrangian departure cells: $c_i
 ightarrow \hat{c}_i$
- Volume correction: $\hat{c}_i \rightarrow \tilde{c}_i$
- Remap from fixed grid to departure grid: $ho_i o ilde
 ho_i, \, ilde m_i = ilde
 ho_i ilde c_i$
- Lagrangian update: $m_i(t+\Delta t) = \tilde{m}_i, \
 ho_i(t+\Delta t) = m_i(t+\Delta t)/c_i$

Application to Semi-Lagrangian Transport

Want to solve:
$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = \mathbf{0}$$

Given cell volume $c_i = \int_{\kappa_i} dV$, cell mass $m_i = \int_{\kappa_i} \rho(\mathbf{x}, t) dV$, and cell average density $\rho_i = m_i/c_i$ at time t



]) Define Lagrangian departure cells: $c_i
ightarrow \hat{c}_i$

- Volume correction: $\hat{c}_i \rightarrow \tilde{c}_i$
- Remap from fixed grid to departure grid: $ho_i o ilde
 ho_i, \ ilde m_i = ilde
 ho_i ilde c_i$
- Lagrangian update: $m_i(t + \Delta t) = \tilde{m}_i, \ \rho_i(t + \Delta t) = m_i(t + \Delta t)/c_i$

We use linear reconstruction of density with Van Leer limiting for remap.

Dukowicz and Baumgardner (2000), Incremental remapping as a transport/advection algorithm, JCP.



Semi-Lagrangian Transport Results

Constant density, rotational flow



Forward Euler with $\Delta t = 0.006$.



Semi-Lagrangian Transport Results

Cylindrical density, rotational flow



Forward Euler with $\Delta t = 0.006$.



Multi-Material Semi-Lagrangian Transport



$$\frac{\partial T_s}{\partial t} + \boldsymbol{u} \cdot \nabla T_s = \boldsymbol{0}, \quad s = 1, \dots, \boldsymbol{3},$$

$$T_{s,i}(t) = \frac{\int_{\kappa_i(t)} T_s \, dV}{\int_{\kappa_i(t)} dV} = \frac{|\kappa_{s,i}(t)|}{|\kappa_i(t)|}$$





Without volume correction:

$$|\Omega_s^{n+1}| = \sum_{i=1}^m T_{s,i}^{n+1} |\kappa_i| = \sum_{i=1}^m \widetilde{T}_{s,i}^n |\kappa_i| \neq \sum_{i=1}^m \widetilde{T}_{s,i}^n |\widetilde{\kappa}_i(t^n)| = |\widetilde{\Omega}_s^n| = |\Omega_s^n|.$$



Conclusions

Presented a new approach for improving the accuracy and physical fidelity of numerical schemes that rely on Lagrangian mesh motion

- Optimization-based volume correction
 - Is computationally efficient
 - Provides significant geometric improvements in corrected meshes
 - Enables semi-Lagrangian transport methods to preserve volumes and constant densities
- Future work
 - Further development of mesh quality constraints and rigorous enforcement of mesh validity
 - Investigate utility of algorithm for mesh quality improvement